

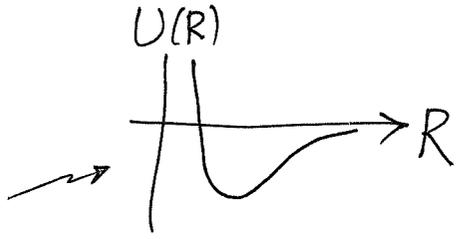
D. Born-Oppenheimer Approximation

[Max Born] [J. Robert Oppenheimer]

- Made approximations to justify separation (decoupling) of Ψ_{molecule} into electronic part and nuclear part

$$\Psi_{\text{molecule}}(\vec{R}; \vec{r}_1, \dots, \vec{r}_N) \approx \Psi_N(\vec{R}) \cdot \Psi_{\text{el}}(\vec{r}_1, \vec{r}_2, \dots, \vec{r}_N; R) \quad (7)$$

nuclei fixed
with separation R

- Why is it a good approximation?
- How bonding comes out from solving Ψ_{el} ? 
- How vibrational and rotational motions come out from $\Psi_N(\vec{R})$?

Meaning:

nuclei k.e. (for electronic problem, ignore for now)

$$\hat{H}_{total} \approx \left(-\frac{\hbar^2}{2\mu} \nabla_{\vec{R}}^2 \right) + \left[\sum_{i=1}^N \left(\frac{-\hbar^2}{2m} \nabla_{\vec{r}_i}^2 \right) - \sum_{i=1}^N \left(\frac{Z_A e^2}{4\pi\epsilon_0 |\vec{r}_i - \vec{R}_A|} + \frac{Z_B e^2}{4\pi\epsilon_0 |\vec{r}_i - \vec{R}_B|} \right) + \frac{Z_A Z_B e^2}{4\pi\epsilon_0 R} + \sum_{i < j} \frac{e^2}{4\pi\epsilon_0 |\vec{r}_i - \vec{r}_j|} \right]$$

\vec{R}_A, \vec{R}_B are regarded fixed $R = |\vec{R}_A - \vec{R}_B|$

a number for fixed R

$$= \left(-\frac{\hbar^2}{2\mu} \nabla_{\vec{R}}^2 \right) + \left[\hat{H}_{electronic}(\vec{p}_1, \dots, \vec{p}_N; \vec{r}_1, \dots, \vec{r}_N; R) \right] \quad (9)$$

involves electrons \rightarrow $\{\hat{p}_i, \hat{r}_i\}$ only

consider this only after solving \hat{H}_{el} Schrodinger problem

involves only electronic variables (with R being a constant)

[Important: Will be back to Eq. (8) later]

Aside: For those who like math...

$$\hat{H} \approx \underbrace{-\frac{\hbar^2}{2\mu} \nabla_{\vec{R}}^2}_{\text{nuclear separation}} + \hat{H}_{\text{electronic}}(\underbrace{\vec{p}_1, \dots, \vec{p}_N; \vec{r}_1, \dots, \vec{r}_N; R}_{\text{electrons' variables only}})$$

[downgraded from operator]
treated as a constant
for now

Nuclear Variables Only

electrons' variables only

Separation of Variables using Eq. (7)

$$\Psi_{\text{molecule}}(\vec{R}; \vec{r}_1, \dots, \vec{r}_N) = \Psi_N(\vec{R}) \cdot \Psi_{\text{el}}(\vec{r}_1, \vec{r}_2, \dots, \vec{r}_N; R)$$

$\Rightarrow \left\{ \begin{array}{l} \text{an equation for } \Psi_{\text{el}}(\vec{r}_1, \vec{r}_2, \dots, \vec{r}_N; R) \text{ (bonding)} \\ \text{an equation for } \Psi_N(\vec{R}) \text{ (vibrational, rotational effects)} \end{array} \right.$

2-step process to solve Eq. (9)

Step 1: Solve electronic part (with R as parameter)

$$\hat{H}_{\text{electronic}} \psi_{\text{el}}(\vec{r}_1, \dots, \vec{r}_N; R) = E_{\text{electronic}}(R) \psi_{\text{el}}(\vec{r}_1, \dots, \vec{r}_N; R) \quad (10)$$

- Many-electron QM problem (c.f. multi-electron atoms)

• Which R to use? [Try different R 's. One R is one problem to solve!]



Eq. (10) gives $E_{\text{el}}(R_1)$



Eq. (10) gives $E_{\text{el}}(R_2)$



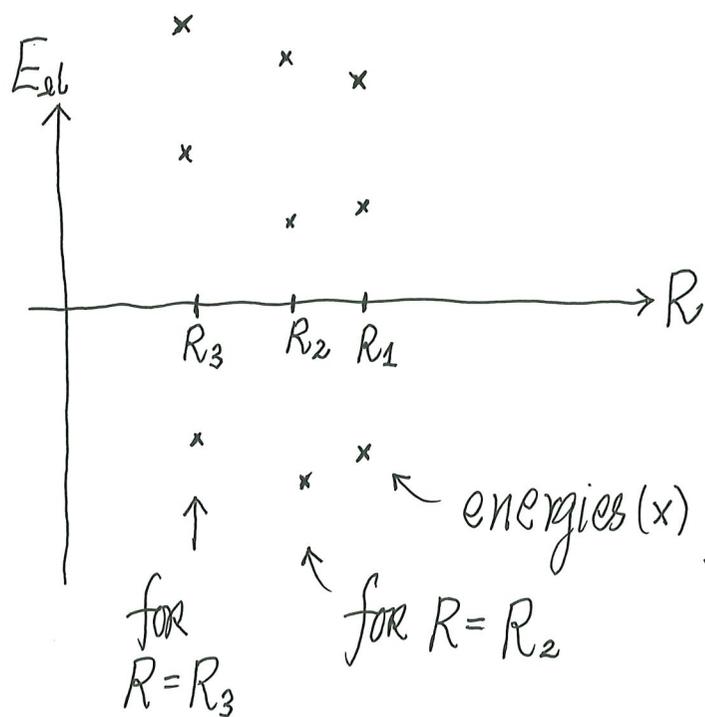
Eq. (10) gives $E_{\text{el}}(R_3)$

and
so
on

⋮

Each is a different
QM problem for
solving Eq. (10)

Key
idea



Labels value of R used in solving TISE[†]

$$\hat{H}_{\text{electronic}} \Psi_{\text{el}}(\vec{r}_1, \dots, \vec{r}_N; R) = E_{\text{el}} \Psi_{\text{el}}(\vec{r}_1, \dots, \vec{r}_N; R) \quad (10)$$

energies (x) solving Eq.(10) for $R=R_1$
 for $R=R_2$
 for $R=R_3$

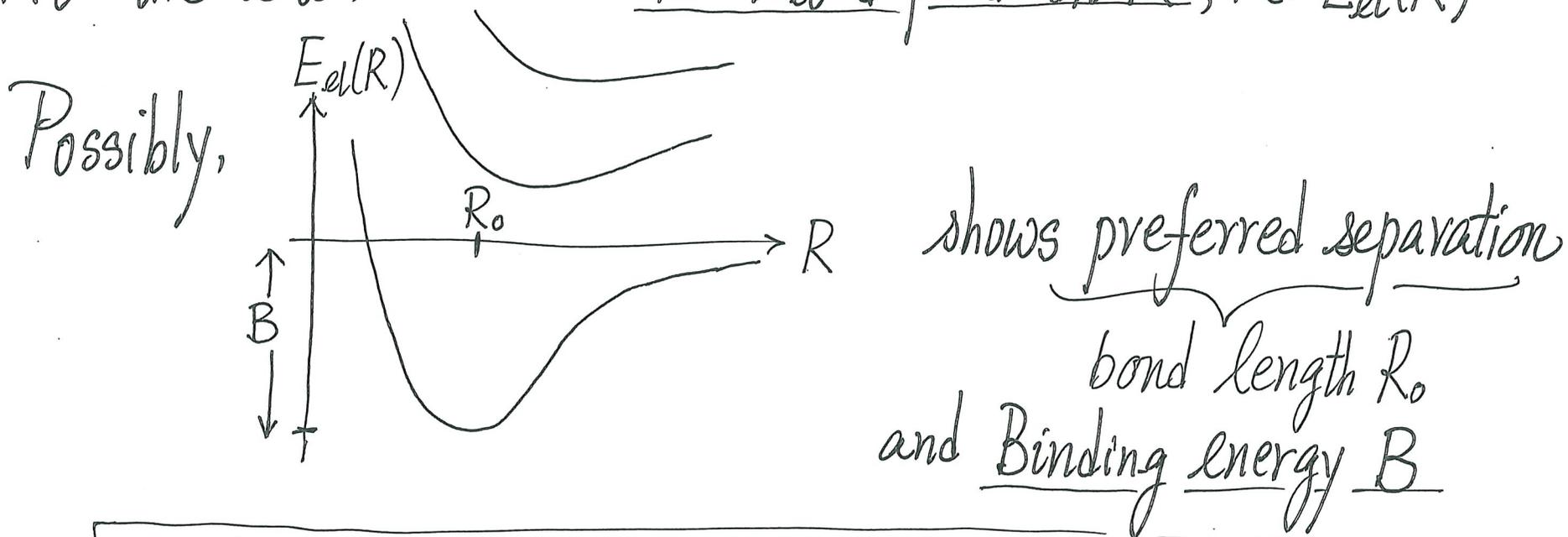
Recall: Solving TISE gives many energy eigenvalues

↑ ↑ ↑
 Different $\hat{H}_{\text{electronic}} \Rightarrow$ Different TISE's

Repeat calculations for many different inter-nucleus separations R
 \Rightarrow many curves $E_{\text{el}}(R)$

[†] Eq.(10) is a many-electron TISE just like atomic problems. How can we "solve" it?
 As in atomic problems, we need to make further approximations.

- At the end, we have how E_{el} depends on R , i.e. $E_{el}(R)$



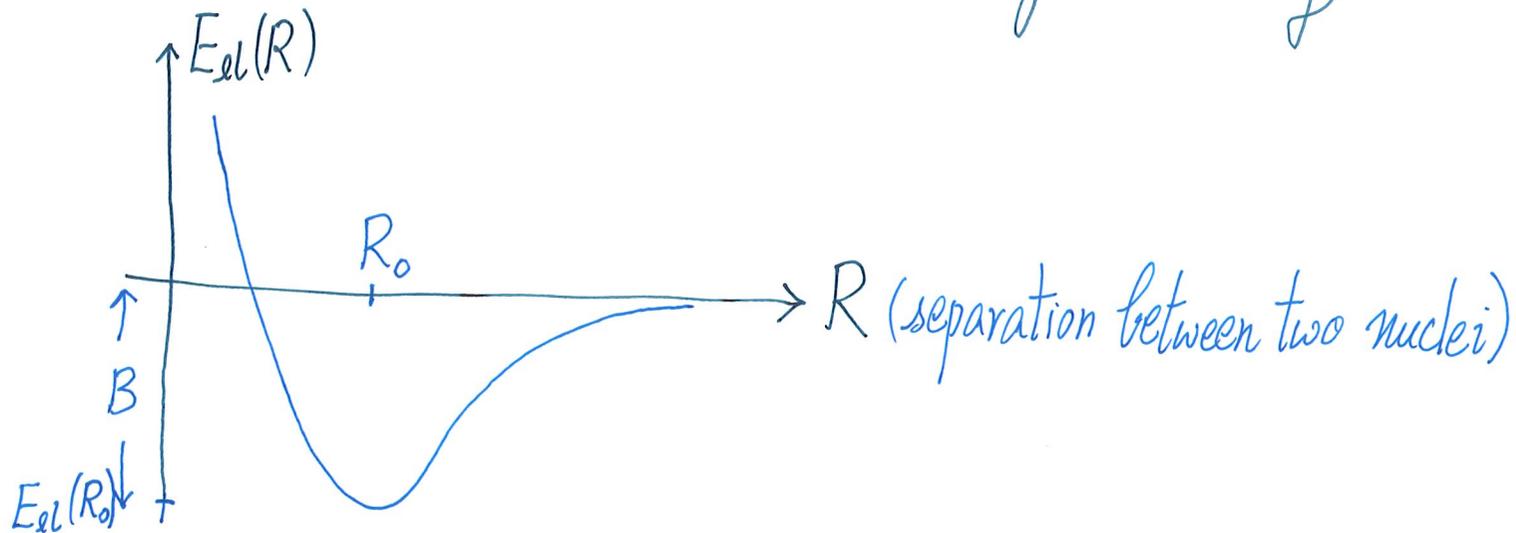
∴ **Electronic problem gives the Quantum Physics of Chemical Bonds!**

Q: Why does the CO (carbon monoxide) molecule have bond length 1.128 \AA ?

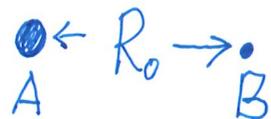
We have $\hat{H}_{\text{electronic}}^{(\text{CO})} \rightarrow \text{TISE} \rightarrow E_{el}(R) \rightarrow R_0$

electrons behavior ($|\Psi_{el}|^2$) that selects preferred separation R_0

Due to the electrons, there emerges binding



If nuclear motion is ignored, electrons [their distribution] select a certain bond length R_0 and certain bond strength B



and energy of the molecule would be $E_{el}(R_0)$

[not true]

The picture of the molecule up to here is

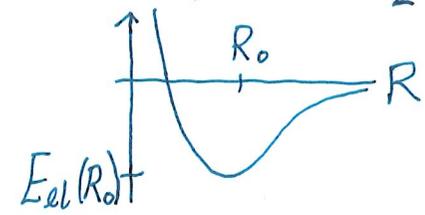


[no vibration, no rotation YET!]

Step 2 : Go back to \hat{H}_{total} and consider nuclear motion
 with $-\frac{\hbar^2}{2\mu} \nabla_R^2$ terms new physics in molecules

At this point, the electrons are done with their duties.

• What's left is the relative motion problem.

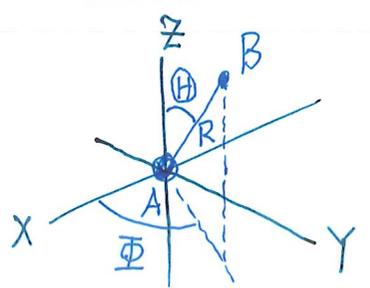


3D One-body of mass μ under the influence of $E_{el}(R)$

depends on separation R only (spherically symmetric)
 [Generally R, Θ, Φ of mass μ]

capital symbols used for nucleus B's coordinates

Picture

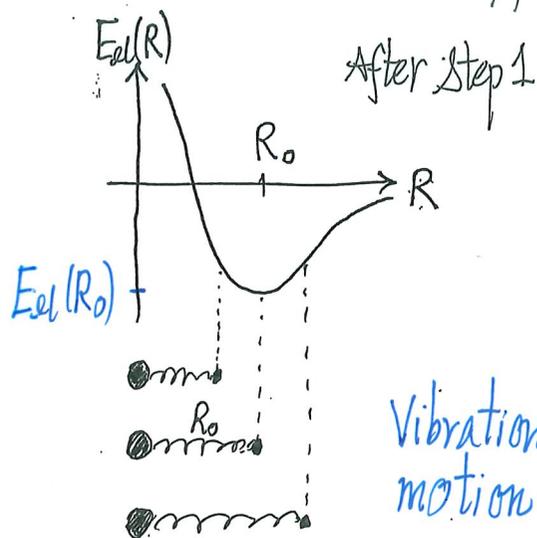
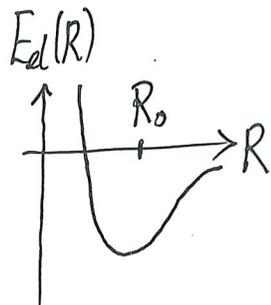


• rod could rotate

• could vibrate (radial direction)

The situation has a 2D analogy (easier to visualize)

Radial motion is approximately harmonic oscillator



Near minimum,

$$E_{el}(R \approx R_0) = E_{el}(R_0) + \frac{1}{2} \left(\frac{d^2 E_{el}}{dR^2} \right)_{E_0} (R - R_0)^2$$

$$= E_{el}(R_0) + \frac{1}{2} k (R - R_0)^2$$

$$k = \frac{1}{2} \left(\frac{d^2 E_{el}}{dR^2} \right)_{E_0} \quad (\text{due to electrons})$$

↑ stiffness of the chemical bond

QM of harmonic oscillator

allowed energies are

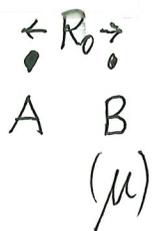
$$(n + \frac{1}{2}) \hbar \omega$$

$$= (n + \frac{1}{2}) \hbar \sqrt{\frac{k}{\mu}} \quad (11)$$

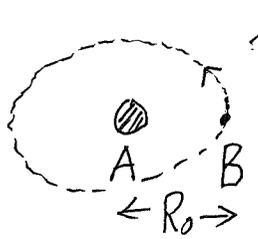
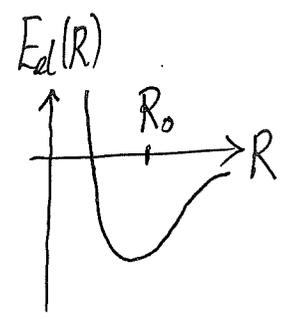
R_0 = equilibrium separation (bond length)
 = "natural length" of spring (bond)

"Spring Constant" ?

▪ Curvature near R_0



2D analogy



rotational motion at the minimum (assume $R=R_0$)

Angular Motion

rotational motion
"2D rigid rotor"

means $R=R_0$

Allowed energies are

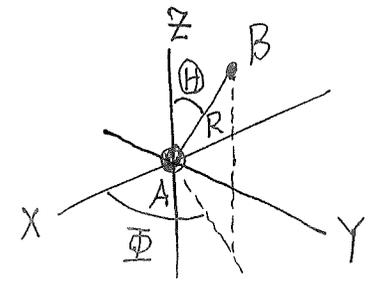
$$\frac{L^2}{2I} = \frac{L^2}{2\mu R_0^2}$$

moment of inertia

Diatomic Molecules

Angular motion is rotational motion

"3D rigid rotor"



B freely moves on sphere of radius R_0

Allowed energies are

$$\frac{l(l+1)\hbar^2}{2\mu R_0^2} \quad (12)^+$$

"rotational levels"

[Chemists use $\frac{J(J+1)\hbar^2}{2\mu R_0^2}$]

At the end (Steps 1+2), an allowed energy of a molecule is specified by

$$E_{\text{total}}^{(\text{molecule})} = \underbrace{E_{\text{el}}(R_0)}_{\text{electronic part}} + \underbrace{\left(n + \frac{1}{2}\right) \hbar \omega}_{\text{vibrational part}} + \underbrace{\frac{l(l+1) \hbar^2}{2\mu R_0^2}}_{\text{rotational part}} \quad (13)$$

which vibrational state
which rotational state

[Key result without going through the mathematics!]

↑
①

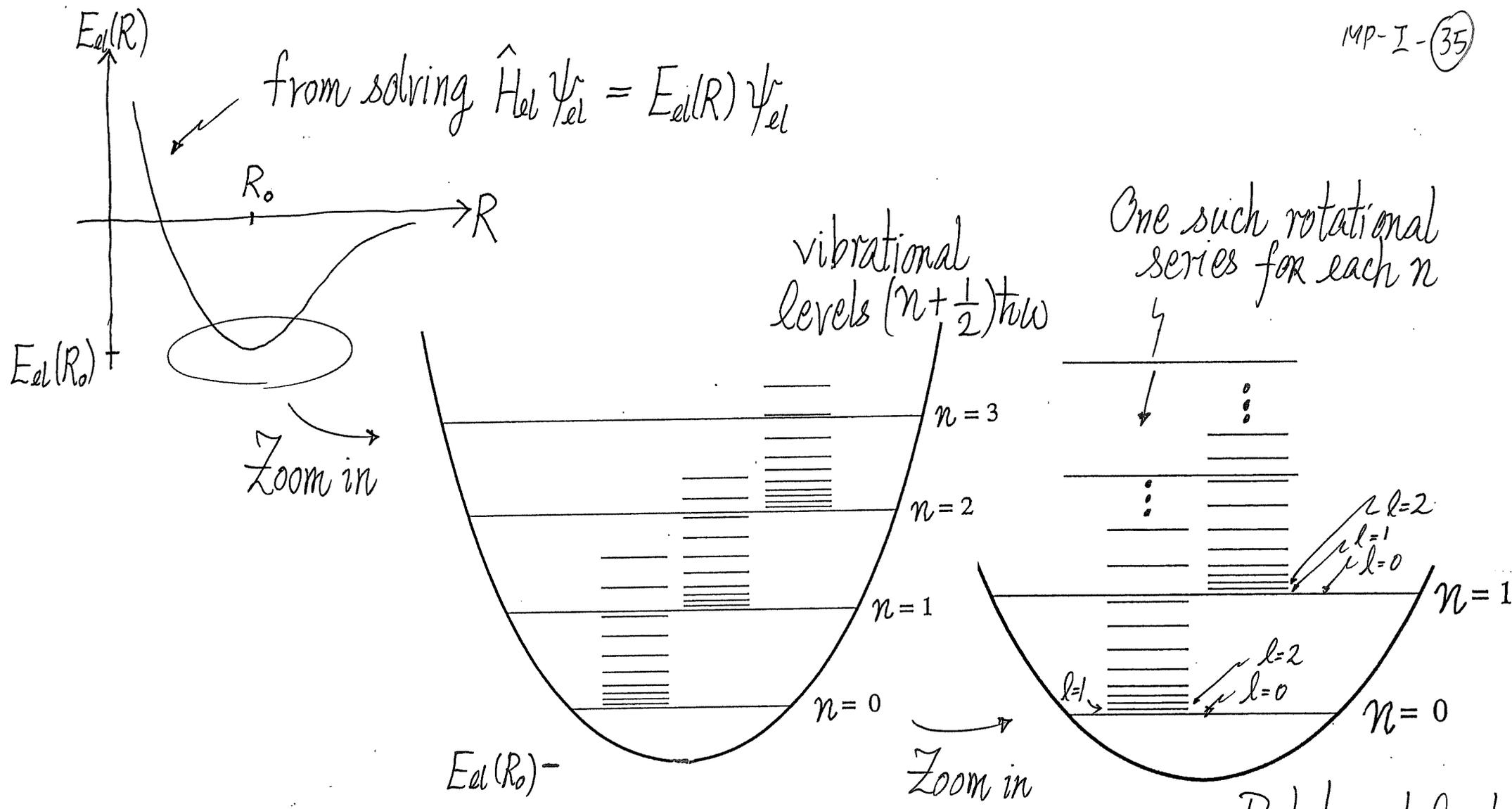
electrons bind the atoms at set the scale of the binding $E_{\text{el}}(R_0)$, the spring constant k (thus ω) and equilibrium bond length R_0 for...

↑
②

... the vibrational motion of the nuclei with $\omega = \sqrt{\frac{k}{\mu}}$...

↑
③

... and the rotational motion of the nuclei



An energy diagram showing the rotational levels associated with each vibrational state in diatomic molecule.

Picture of Key Result Eq. (13)

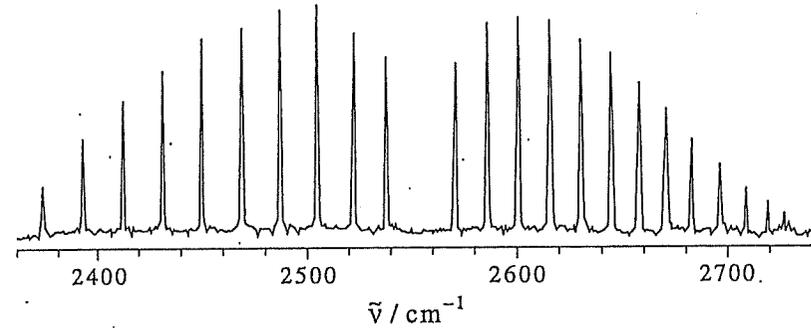
$$\frac{l(l+1)\hbar^2}{2I}$$

[each l , $(2l+1)$ values of m]
degeneracy

Experimental consequences

- Transitions between one molecular state $[n, l]$ to another state $[n', l']$ give finger-print molecular spectrum of the molecule

IR range



Rotational-vibrational spectrum of HBr
 [Questions: Get R_0 and $\omega = \sqrt{\frac{k}{\mu}}$ out from spectrum?]

- Must understand the physical picture behind Born-Oppenheimer approximation first, before going into more detail
- It is important also because the same physics works in understanding solids. ["Solid is just a huge molecule"]
- Eq.(13) is the key result for understanding molecular spectrum
[c.f. atomic spectra motivated the development of QM]
- This ends the discussion on the Big Picture of the QM theory of molecular states